

Winter-Semester 2017/18

Moderne Theoretische Physik IIIa
Statistische Physik

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Institut für Theorie der Kondensierten Materie

Do 11:30-13:00, Lehmann Raum 022, Geb 30.22

<http://www.tkm.kit.edu/lehre/>

Ideales Fermi-Gas, T=0

$$n = \frac{N}{V} = (2s + 1) \int d\epsilon \nu(\epsilon) \theta(\mu - \epsilon)$$

Die Teilchen-Dichte n ist gegeben $\rightarrow \mu(T = 0) \equiv \epsilon_F$

ϵ_F - Fermi-Energie

$$p_F = \sqrt{2m\epsilon_F}$$

$$\nu(\epsilon) = \frac{m^{3/2}\epsilon^{1/2}}{\sqrt{2}\pi^2\hbar^3}$$

Zustandsdichte

$$\mu(T = 0) = \epsilon_F = \frac{\hbar^2}{2m} \left(\frac{6\pi^2 n}{2s + 1} \right)^{\frac{2}{3}}$$

Ideales Fermi-Gas

Entartetes Fermi-Gas $k_{\text{B}}T \ll \epsilon_F$

$$Z_G(T, V, \mu) = \prod_{\lambda} \left[1 + e^{-\beta(\epsilon_{\lambda} - \mu)} \right]$$

$$\Omega(T, V, \mu) = -k_{\text{B}}T \ln Z_G(T, V, \mu) = -k_{\text{B}}T \sum_{\lambda} \ln \left[1 + e^{-\beta(\epsilon_{\lambda} - \mu)} \right]$$

$$\Omega(T, V, \mu) = -k_{\text{B}}T(2s+1)V \int_{-\infty}^{\infty} d\epsilon \nu(\epsilon) \ln \left[1 + e^{-\beta(\epsilon - \mu)} \right]$$

2 x partielle Integration

$$\Omega(T, V, \mu) = -(2s+1)V \int_{-\infty}^{\infty} d\epsilon a(\epsilon) n_F(\epsilon)$$

$$a(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 \nu(\epsilon_1)$$

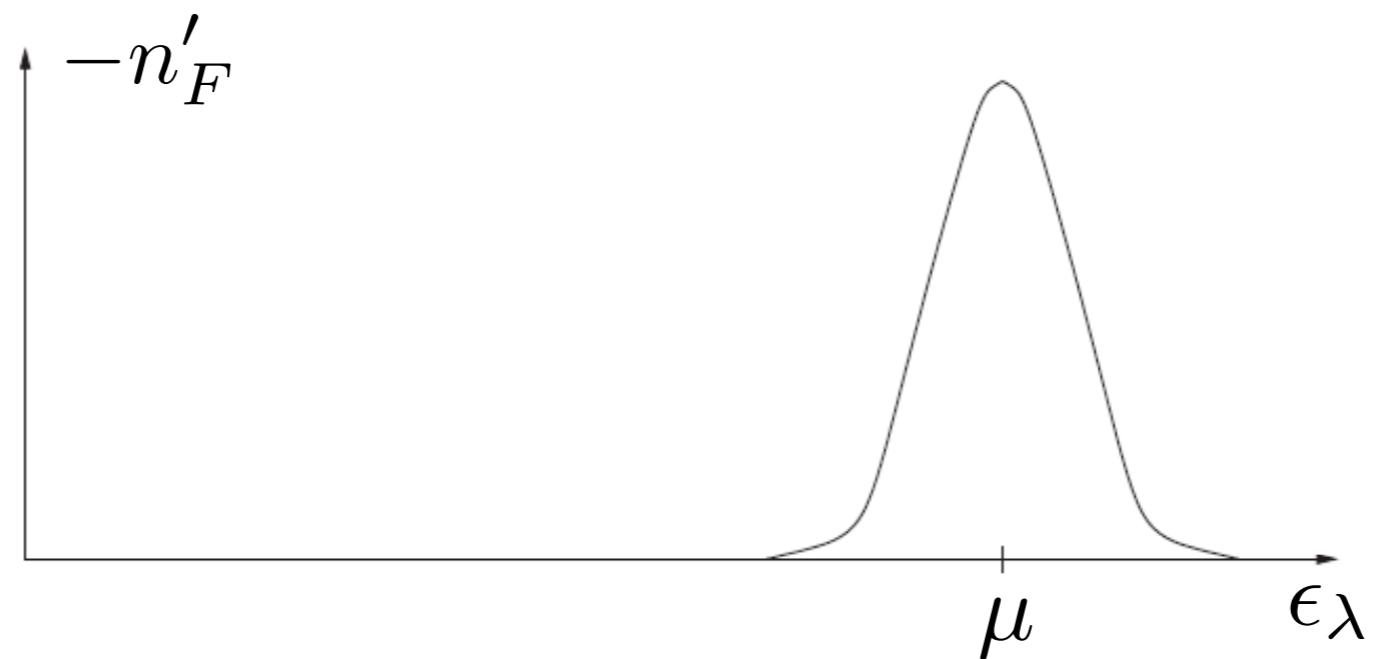
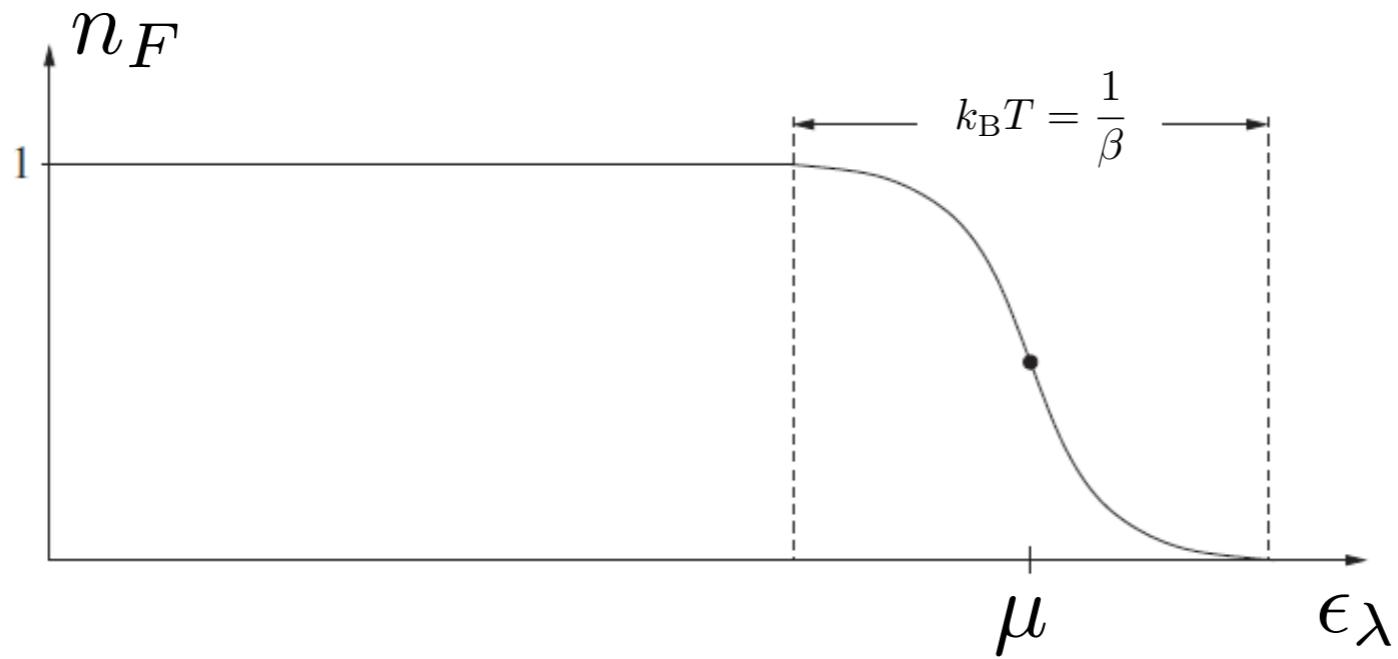
$$\Omega(T, V, \mu) = -(2s+1)V \int_{-\infty}^{\infty} d\epsilon b(\epsilon) (-n'_F(\epsilon))$$

$$b(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 a(\epsilon_1)$$

Sommerfeld-Entwicklung

$$k_B T \ll \epsilon_F$$

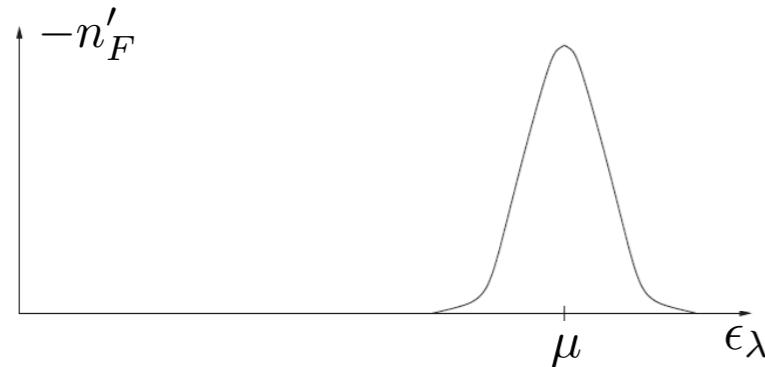
$$-\frac{\partial n_F}{\partial \epsilon_\lambda} = \frac{1}{4k_B T \cosh^2 \frac{\epsilon_\lambda - \mu}{2k_B T}}$$



Ideales Fermi-Gas

Entartetes Fermi-Gas $k_{\text{B}}T \ll \epsilon_F$

$$\Omega(T, V, \mu) = -(2s+1)V \int_{-\infty}^{\infty} d\epsilon b(\epsilon)(-n'_F(\epsilon))$$



$$a(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 \nu(\epsilon_1)$$

$$b(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 a(\epsilon_1)$$

für $\epsilon \approx \mu$ gilt $b(\epsilon) = b(\mu) + a(\mu)(\epsilon - \mu) + \frac{1}{2} \nu(\mu)(\epsilon - \mu)^2 + \dots$

$$\int_{-\infty}^{\infty} dx \frac{x^2}{\cosh^2 x} = \frac{\pi^2}{6}$$

$$\Omega(T, V, \mu) = -(2s+1)V \left[b(\mu) + \frac{\pi^2}{6} (k_{\text{B}}T)^2 \nu(\mu) + \dots \right]$$

Sommerfeld-Entwicklung

Ideales Fermi-Gas

$$a(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 \nu(\epsilon_1)$$

Entartetes Fermi-Gas $k_{\text{B}}T \ll \epsilon_F$

$$b(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 a(\epsilon_1)$$

$$\Omega(T, V, \mu) = -(2s+1)V \left[b(\mu) + \frac{\pi^2}{6} (k_{\text{B}}T)^2 \nu(\mu) + \dots \right]$$

$$N = -\frac{\partial \Omega}{\partial \mu} \Big|_{T,V} = (2s+1)V \left[a(\mu) + \frac{\pi^2}{6} (k_{\text{B}}T)^2 \frac{d\nu(\mu)}{d\mu} + \dots \right]$$

$$a(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 \nu(\epsilon_1) = \frac{N(T=0, V, \mu)}{(2s+1)V}$$

$$N(T, V, \mu) = N(T=0, V, \mu) + (2s+1)V \frac{\pi^2}{6} (k_{\text{B}}T)^2 \frac{d\nu(\mu)}{d\mu} + \dots$$

$$n(T, \mu) = n(T=0, \mu) + (2s+1) \frac{\pi^2}{6} (k_{\text{B}}T)^2 \frac{d\nu(\mu)}{d\mu} + \dots$$

Entartetes ideales Fermi-Gas $k_{\text{B}}T \ll \epsilon_F$

Das chemische Potential μ ist gegeben $\rightarrow n(T, \mu) = ?$

$$n(T, \mu) = n(T=0, \mu) + (2s+1) \frac{\pi^2}{6} (k_{\text{B}}T)^2 \frac{d\nu(\mu)}{d\mu} + \dots$$

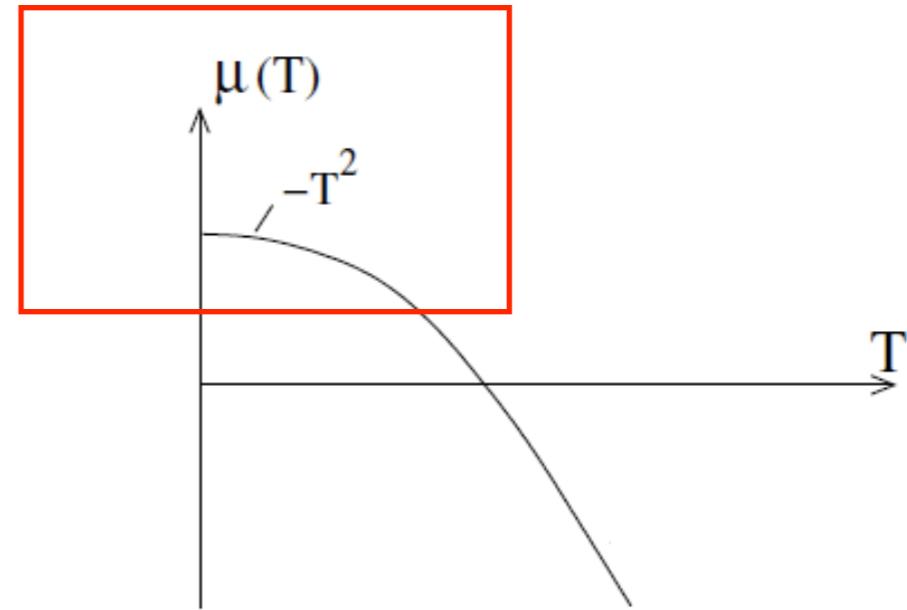
$$\mu(T=0) = \epsilon_F = \frac{\hbar^2}{2m} \left(\frac{6\pi^2 n}{2s+1} \right)^{\frac{2}{3}} \quad \rightarrow \quad n(T=0, \mu) = \frac{2s+1}{6\pi^2} \left(\frac{2m\mu}{\hbar^2} \right)^{\frac{3}{2}}$$

Die Teilchen-Dichte n ist gegeben $\rightarrow \mu(T, n) = ?$ $\nu(\epsilon) = \frac{m^{3/2}\epsilon^{1/2}}{\sqrt{2\pi^2\hbar^3}}$

$$n = \frac{2s+1}{6\pi^2} \left(\frac{2m\epsilon_F}{\hbar^2} \right)^{\frac{3}{2}} = \frac{2s+1}{6\pi^2} \left(\frac{2m}{\hbar^2} \right)^{\frac{3}{2}} \left[\mu^{3/2} + \frac{\pi^2}{8\sqrt{\mu}} (k_{\text{B}}T)^2 + \dots \right]$$

$$\epsilon_F^{3/2} = \mu^{3/2} + \frac{\pi^2}{8\sqrt{\mu}} (k_{\text{B}}T)^2 + \dots$$

$$\mu = \epsilon_F \left[1 - \frac{\pi^2}{12} \left(\frac{k_{\text{B}}T}{\epsilon_F} \right)^2 + \dots \right]$$



Ideales Fermi-Gas

$$a(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 \nu(\epsilon_1)$$

Entartetes Fermi-Gas $k_B T \ll \epsilon_F$

$$b(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 a(\epsilon_1)$$

$$\Omega(T, V, \mu) = -(2s+1)V \left[b(\mu) + \frac{\pi^2}{6} (k_B T)^2 \nu(\mu) + \dots \right]$$

$$S(T, V, \mu) = -\frac{\partial \Omega}{\partial T} \Big|_{V, \mu} = (2s+1)V \frac{\pi^2 \nu(\mu)}{3} k_B^2 T$$

kanonisch - großkanonisch: **Wachsamkeit nötig** $\mu = \mu(T, N)$

$$C_V = T \left(\frac{\partial S}{\partial T} \right)_{V, N} = T \left(\frac{\partial S(T, V, \mu(N, T))}{\partial T} \right)_{V, N}$$

$$C_V = \textcolor{red}{T} \left(\frac{\partial S(T, V, \mu)}{\partial T} \right)_{V, \mu} + \textcolor{blue}{T} \left(\frac{\partial S(T, V, \mu)}{\partial \mu} \right)_{V, T} \left(\frac{\partial \mu(T, N)}{\partial T} \right)_N$$

$$C_V \approx N k_B \frac{\pi^2}{2} \frac{k_B T}{\epsilon_F} + O(T^3)$$

Ideales Fermi-Gas

Entartetes Fermi-Gas $k_{\text{B}}T \ll \epsilon_F$

$$\Omega(T, V, \mu) = -(2s+1)V \left[b(\mu) + \frac{\pi^2}{6} (k_{\text{B}}T)^2 \nu(\mu) + \dots \right]$$

$$P = -\frac{\partial \Omega}{\partial V} \Big|_{T, \mu} = (2s+1) \left[b(\mu) + \frac{\pi^2}{6} (k_{\text{B}}T)^2 \nu(\mu) + \dots \right]$$

$$P = P(T=0, V, \mu) + (2s+1) \frac{\pi^2}{6} (k_{\text{B}}T)^2 \nu(\mu) + \dots$$

$$a(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 \nu(\epsilon_1) = \frac{N(T=0, \mu=\epsilon, V)}{(2s+1)V} = \frac{1}{(2\pi\hbar)^3} \cdot \frac{4}{3} \pi (2m\epsilon)^{3/2}$$

$$b(\mu) \equiv \int_{-\infty}^{\mu} d\epsilon_1 a(\epsilon_1) = \frac{2}{5} \mu \cdot \frac{N(T=0, \mu, V)}{(2s+1)V}$$

$$P(T=0, V, \mu)V = (2s+1)Vb(\mu) = \frac{2}{5}N(T=0, V, \mu)\mu = \frac{2}{5}Nk_{\text{B}}T_F$$

Ideales Fermi-Gas

Beliebige Temperaturen

$$\Omega(T, V, \mu) = -k_B T \sum_{\lambda} \ln \left[1 + e^{-\beta(\epsilon_{\lambda} - \mu)} \right]$$

$$\begin{aligned}\Omega(T, V, \mu) &= -k_B T (2s+1) V \int d\epsilon \nu(\epsilon) \ln \left[1 + e^{-\beta(\epsilon - \mu)} \right] \\ &= -(2s+1) \frac{V}{\lambda_T^3} k_B T f_{5/2}(z) \quad z \equiv e^{\beta\mu}\end{aligned}$$

$$\nu(\epsilon) = \frac{m^{3/2} \epsilon^{1/2}}{\sqrt{2\pi^2 \hbar^3}}$$

$$\lambda_T = \sqrt{\frac{2\pi\hbar^2}{mk_B T}}$$

$$f_{5/2}(z) \equiv \frac{4}{\sqrt{\pi}} \int_0^{\infty} dx x^2 \ln(1 + ze^{-x^2}) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1} z^n}{n^{5/2}}$$

$$N = -\frac{\partial \Omega}{\partial \mu} \Big|_{T,V} = (2s+1) \frac{V}{\lambda_T^3} f_{3/2}(z)$$

$$f_{3/2}(z) \equiv \sum_{n=1}^{\infty} \frac{(-1)^{n+1} z^n}{n^{3/2}}$$

Ideales Fermi-Gas

Beliebige Temperaturen

$$z \equiv e^{\beta\mu} \quad \lambda_T = \sqrt{\frac{2\pi\hbar^2}{mk_B T}}$$

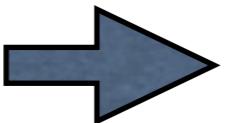
$$\Omega(T, V, \mu) = -(2s+1) \frac{V}{\lambda_T^3} k_B T f_{5/2}(z)$$

$$f_{5/2}(z) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1} z^n}{n^{5/2}}$$

$$N = -\frac{\partial \Omega}{\partial \mu} \Big|_{T,V} = (2s+1) \frac{V}{\lambda_T^3} f_{3/2}(z)$$

$$f_{3/2}(z) \equiv \sum_{n=1}^{\infty} \frac{(-1)^{n+1} z^n}{n^{3/2}}$$

$$S(T, V, \mu) = -\frac{\partial \Omega}{\partial T} \Big|_{V,\mu}$$



$$U = -\frac{3}{2}\Omega = \frac{3}{2}PV \quad \textcolor{blue}{Übung}$$

Ideales Fermi-Gas

$$z \equiv e^{\beta\mu} \quad \lambda_T = \sqrt{\frac{2\pi\hbar^2}{mk_B T}}$$

Hohe Temperaturen $k_B T \gg \epsilon_F$

Vermutung: $z \ll 1$

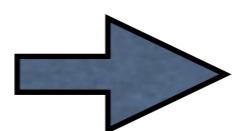
$$N = (2s+1) \frac{V}{\lambda_T^3} f_{3/2}(z) \approx (2s+1) \frac{V}{\lambda_T^3} z$$

$$f_{3/2}(z) \equiv \sum_{n=1}^{\infty} \frac{(-1)^{n+1} z^n}{n^{3/2}}$$

$$\Omega(T, V, \mu) = -(2s+1) \frac{V}{\lambda_T^3} k_B T f_{5/2}(z) \approx -k_B T N$$

$$f_{5/2}(z) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1} z^n}{n^{5/2}}$$

$$U = -\frac{3}{2}\Omega = \frac{3}{2}PV$$



$$U = \frac{3}{2} k_B T N$$
$$PV = k_B T N$$

ideales klassisches Gas

Ideales Fermi-Gas

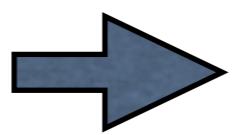
Hohe Temperaturen $k_{\text{B}}T \gg \epsilon_F$

$$z \equiv e^{\beta\mu}$$

$$\lambda_T = \sqrt{\frac{2\pi\hbar^2}{mk_{\text{B}}T}}$$

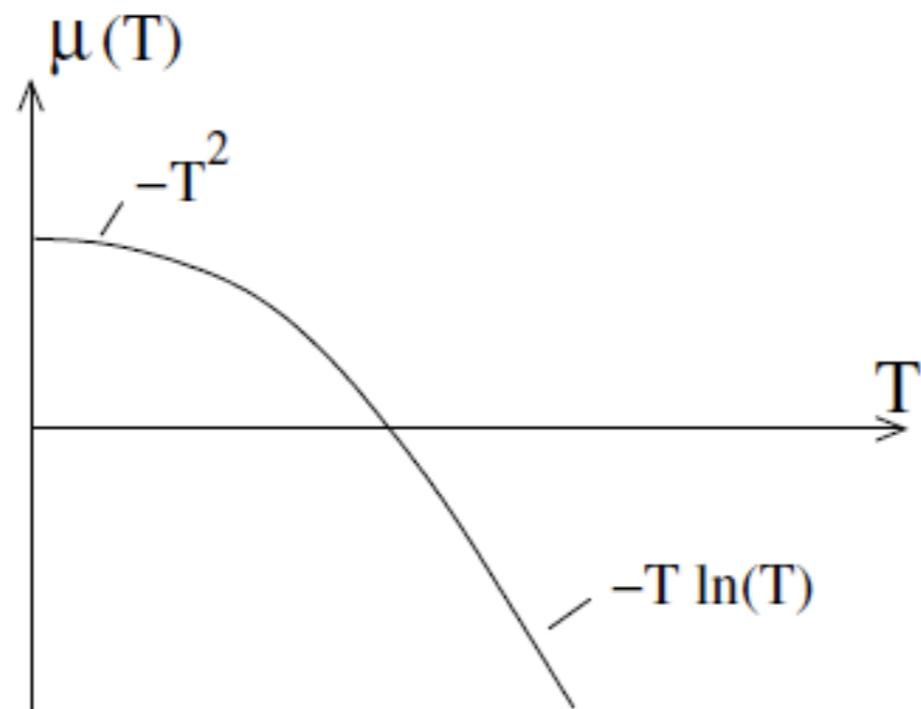
Vermutung: $z \ll 1$

$$N \approx (2s + 1) \frac{V}{\lambda_T^3} z$$



$$\mu \approx k_{\text{B}}T \ln \frac{\lambda_T^3 n}{(2s + 1)}$$

$$n \equiv \frac{N}{V}$$



Ideales Fermi-Gas

$$U = -\frac{3}{2}\Omega = \frac{3}{2}PV$$

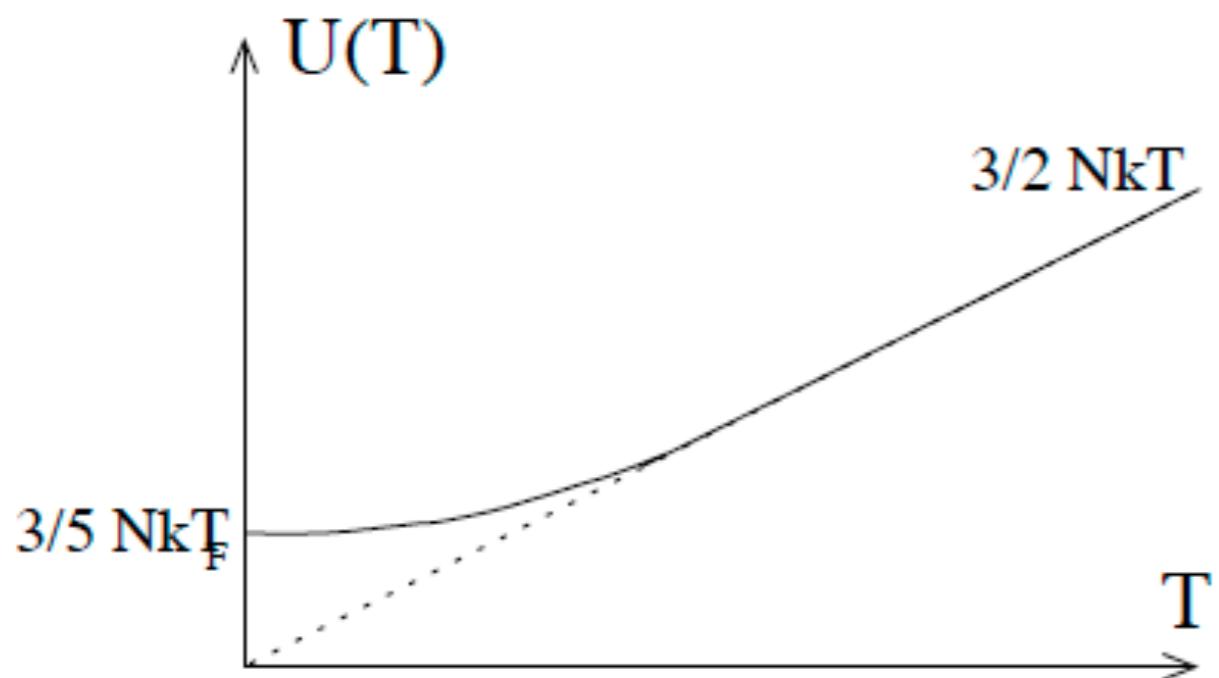
$$P(T=0)V = \frac{2}{5}Nk_B T_F$$

$$k_B T \ll \epsilon_F \rightarrow$$

$$U \approx \frac{3}{5}Nk_B T_F$$

$$k_B T \gg \epsilon_F \rightarrow$$

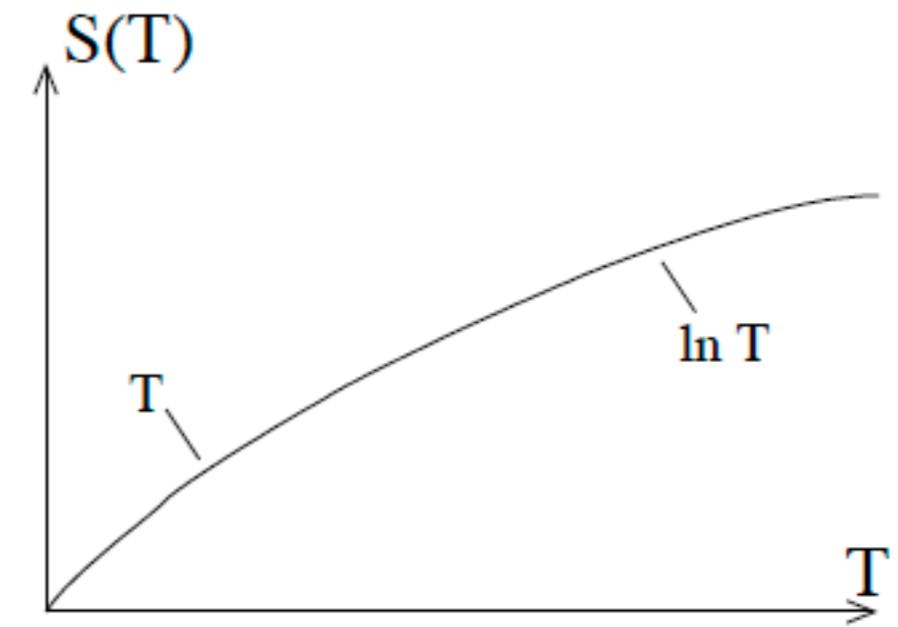
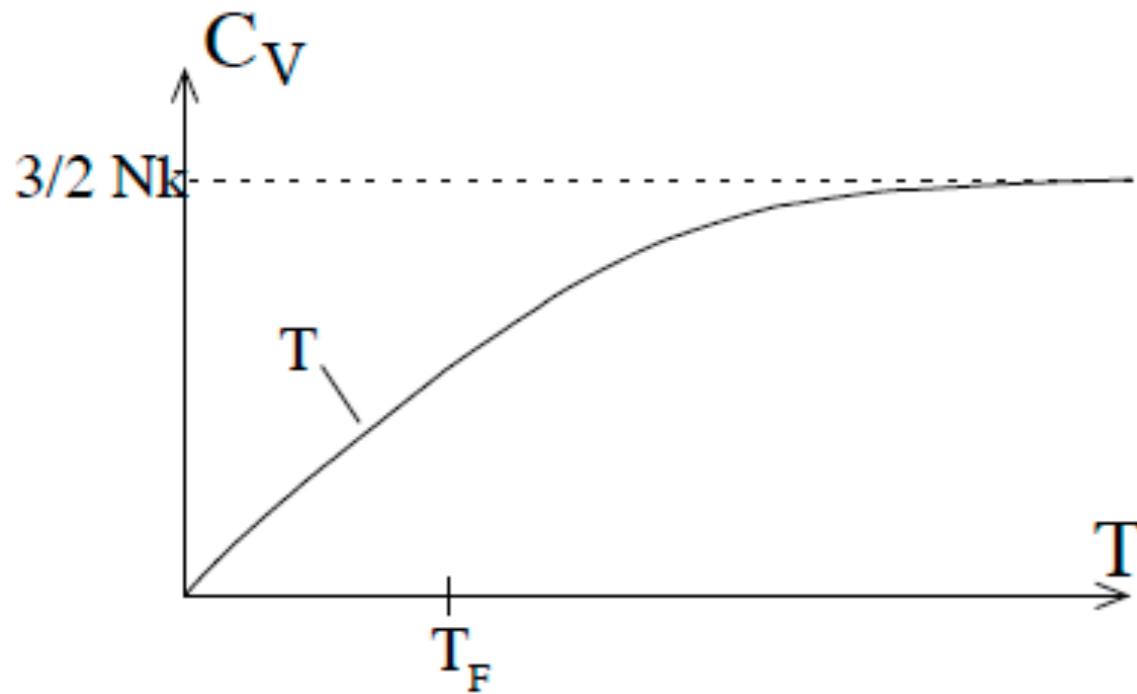
$$U = \frac{3}{2} k_B T N$$



Ideales Fermi-Gas

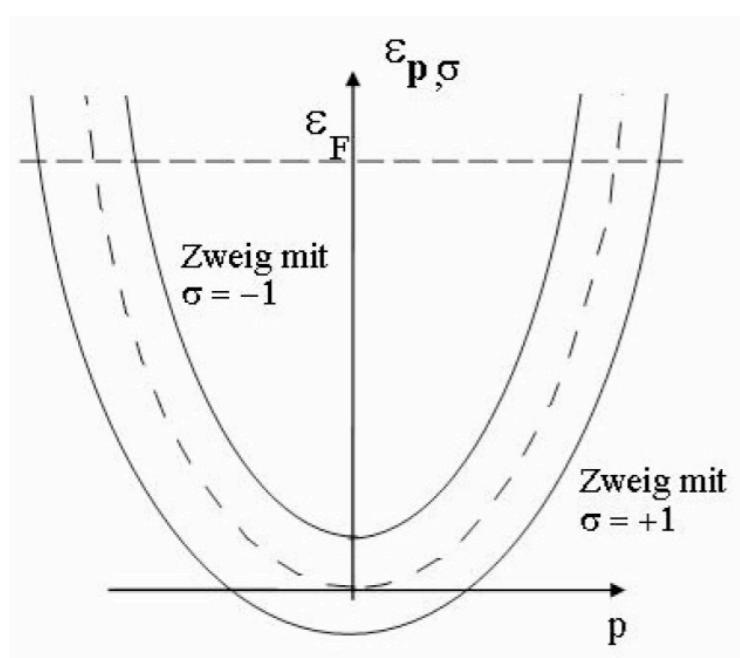
$$k_B T \ll \epsilon_F \rightarrow C_V \approx N k_B \frac{\pi^2}{2} \frac{k_B T}{\epsilon_F}$$

$$k_B T \gg \epsilon_F \rightarrow C_V \approx \frac{3}{2} N k_B$$



Ideales Fermi-Gas

Pauli-Paramagnetismus



$$H = \frac{(\mathbf{p} - \frac{e}{c}\mathbf{A})^2}{2m} - \mu_B \sigma \mathbf{B} \quad \mu_B = \frac{e\hbar}{2mc}$$

Diamagnetismus Paramagnetismus

$$\epsilon_{\mathbf{p},\sigma} = \frac{\mathbf{p}^2}{2m} - \mu_B \sigma B = \epsilon_{\mathbf{p}} - \mu_B \sigma B$$

$$Z_G = \sum_{n_{p,+}, n_{p,-}} e^{-\beta [(\epsilon_{\mathbf{p}} - \mu_B B - \mu)n_{\mathbf{p},+} + (\epsilon_{\mathbf{p}} + \mu_B B - \mu)n_{\mathbf{p},-}]} \quad \mu_{\pm} = \mu \pm \mu_B B$$

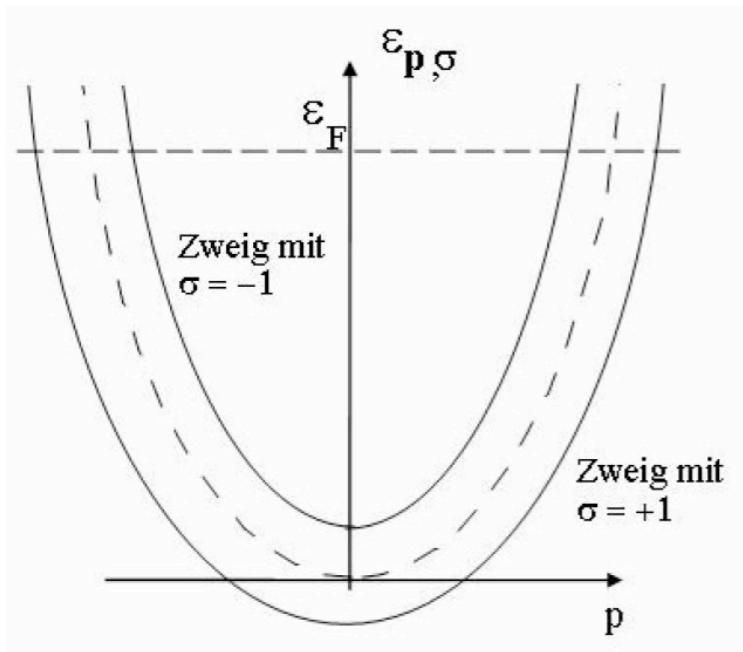
$$Z_G = Z_G^0(\mu_+) Z_G^0(\mu_-) \rightarrow \Omega(T, V, \mu) = \Omega^0(T, V, \mu_+) + \Omega^0(T, V, \mu_-)$$

$$\Omega^0(T, V, \mu) = -k_B T V \int d\epsilon \nu(\epsilon) \ln \left[1 + e^{-\beta(\epsilon - \mu)} \right] = -\frac{V}{\lambda_T^3} k_B T f_{5/2}(z)$$

Ideales Fermi-Gas

Pauli-Paramagnetismus

$$\mu_{\pm} = \mu \pm \mu_B B$$



$$a(\epsilon) \equiv \int_{-\infty}^{\epsilon} d\epsilon_1 \nu(\epsilon_1)$$

$$\nu(\epsilon) = \frac{m^{3/2} \epsilon^{1/2}}{\sqrt{2} \pi^2 \hbar^3}$$

$$\frac{N_+ - N_-}{V} = a(\mu_+) - a(\mu_-) \approx 2\nu(\epsilon_F)\mu_B B$$

$$k_B T \ll \epsilon_F$$

$$\frac{M}{V} = \mu_B \frac{(N_+ - N_-)}{V} = 2\nu(\epsilon_F)\mu_B^2 B = \chi B$$

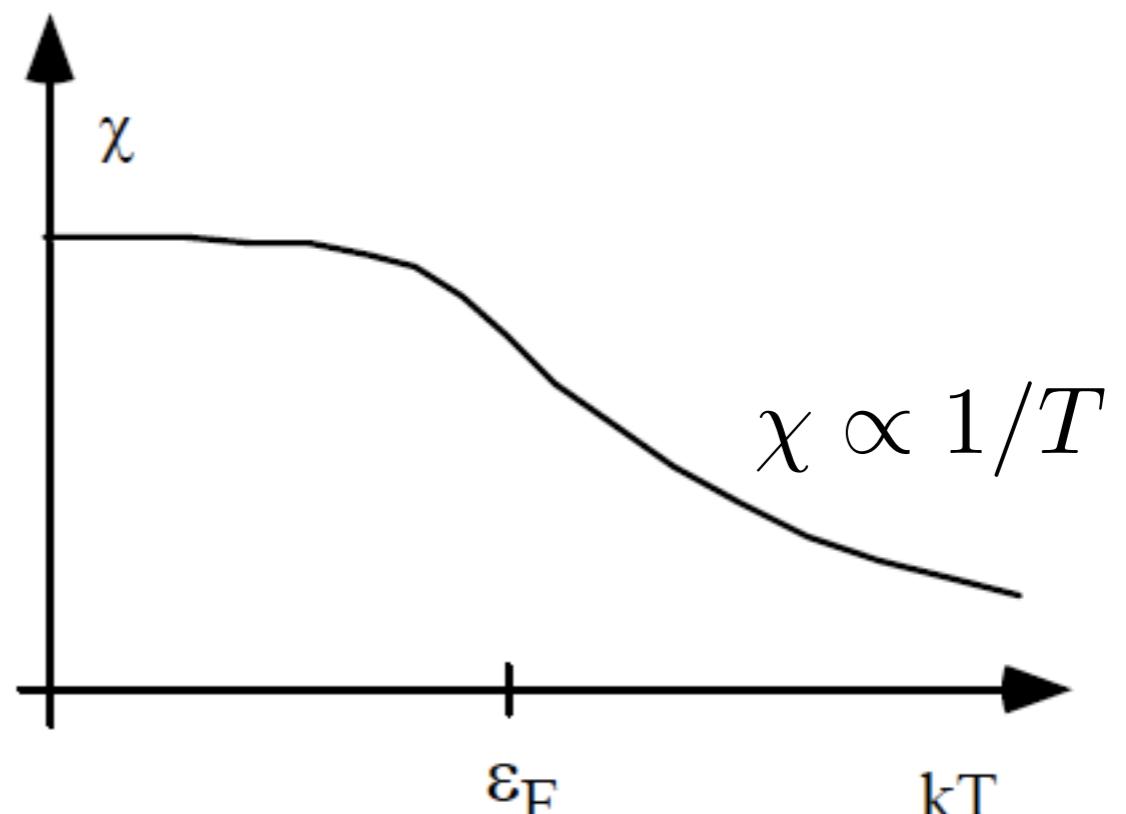
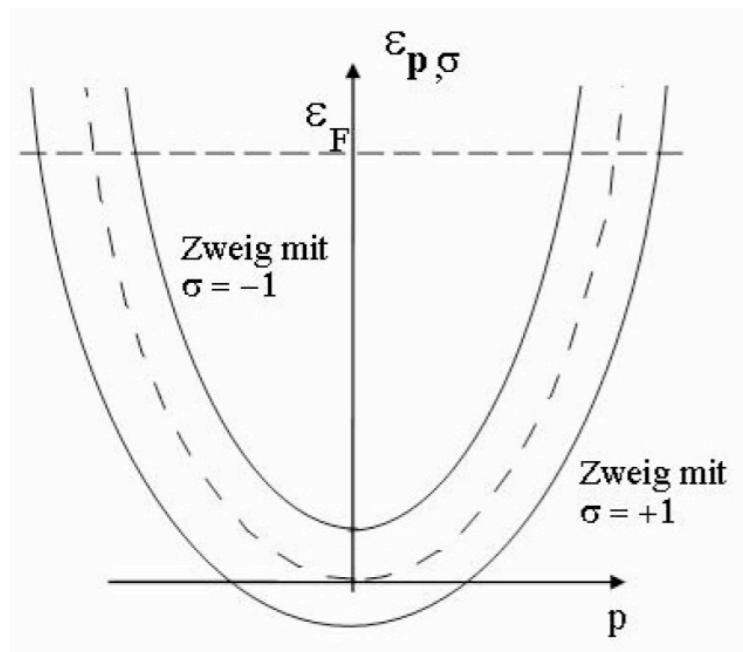
$$k_B T \gg \epsilon_F$$

$$\begin{aligned} \frac{M}{V} &= \mu_B \frac{(N_+ - N_-)}{V} \approx \frac{\mu_B}{\lambda_T^3} (e^{\beta\mu_+} - e^{\beta\mu_-}) \\ &= \frac{\mu_B}{\lambda_T^3} e^{\beta\mu} \beta (2\mu_B B) = n\beta\mu_B^2 B = \chi B \end{aligned}$$

$$N \approx (2s+1) \frac{V}{\lambda_T^3} z$$

Ideales Fermi-Gas

Pauli-Paramagnetismus



$$k_B T \ll \epsilon_F$$

$$\chi = 2\nu(\epsilon_F)\mu_B^2$$

$$\nu(\epsilon) = \frac{m^{3/2}\epsilon^{1/2}}{\sqrt{2\pi^2}\hbar^3}$$

$$k_B T \gg \epsilon_F$$

$$\chi = n\beta\mu_B^2 = \frac{n\mu_B^2}{k_B T}$$