

Feldtheorie der Kondensierten Materie, RG und Quantenkritikalität SS 2013Prof. Dr. J. Schmalian
Dr. P. Orth**Blatt 01**
Besprechung 23.04.2013**1. Ising ϕ^4 -field theory** (10 + 20 + 10 + 10 + 10 + 10 = 70 Punkte)

In this exercise we derive the scalar ϕ^4 -field theory starting from the microscopic Ising Hamiltonian

$$H = -\frac{1}{2} \sum_{i,j=1}^N J_{ij} S_i S_j, \quad (1)$$

where i, j run over lattice sites (*e.g.* of a d -dimensional cubic lattice), the Ising variables take values $S_i = \pm 1$, $J_{ij} > 0$ denotes a ferromagnetic coupling between spins.

(a) Prove the following identity for multi-dimensional integrals over real variables

$$\int_{-\infty}^{\infty} \frac{dx_1 \cdots dx_n}{(2\pi)^{n/2}} e^{-\frac{1}{2} \sum_{i,j} x_i A_{ij} x_j + \sum_i x_i J_i} = [\det A]^{-1/2} e^{\frac{1}{2} \sum_{ij} J_i (A^{-1})_{ij} J_j},$$

where A is a real symmetric positive definite matrix.

(b) Use this identity to make the exponent of the partition sum for the Ising model linear in the spin variables S_i (Hubbard-Stratonovich transformation). The partition sum reads

$$Z = \sum_{\{S_i\}} \exp(-\beta H).$$

The price to pay is to introduce a (Hubbard-Stratonovich) real variable x_i at each lattice site.

(c) It is now possible to exactly perform the summation over Ising variables $\sum_{\{S_i\}}$, which leads to a spatially local (potential) term in x_j . Then perform the variable transformation $\phi_i = \frac{1}{\sqrt{2}} \sum_j K_{ij}^{-1} x_j$ where $K_{ij} = \beta J_{ij}$ and x_j are the fields introduced in part (b).

(d) Now expand the exponent $H_{\text{eff}}[\phi_j]$ in $Z \sim \int \mathcal{D}[\phi_j] \exp(-\beta H_{\text{eff}}[\phi_j])$ up to quartic order in ϕ_j and then go to momentum space. Here, $\mathcal{D}[\phi_j]$ denotes the integration measure. Explicitly calculate the coupling function $J(\mathbf{k})$ in momentum space (assuming translational invariance $J_{ij} = J_{\mathbf{r}=\mathbf{r}_i-\mathbf{r}_j}$)

$$J(\mathbf{k}) = \sum_{\mathbf{r}} J_{\mathbf{r}} e^{i\mathbf{k} \cdot (\mathbf{r}_i - \mathbf{r}_j)}$$

assuming ferromagnetic nearest-neighbor coupling on a d -dimensional cubic lattice.

(e) Expand $J(\mathbf{k})$ up to second order in \mathbf{k}^2 , and identify the temperature T_c at which the (T -dependent) constant in front of the \mathbf{k} -independent part of the quadratic term vanishes and changes sign.

Then expand the coefficient in front of the quadratic \mathbf{k}^2 term as well as the quartic term to lowest order in $T - T_c$ to bring the expression into the form

$$H_{\text{eff}}[\phi_{\mathbf{k}}] = \frac{1}{2} \int \frac{d^d k}{(2\pi)^d} \phi_{\mathbf{k}} [a_0(T - T_c) + b\mathbf{k}^2] \phi_{-\mathbf{k}} + \frac{u}{4} \int \frac{d^d k_1}{(2\pi)^d} \frac{d^d k_2}{(2\pi)^d} \frac{d^d k_3}{(2\pi)^d} \phi_{\mathbf{k}_1} \phi_{\mathbf{k}_2} \phi_{\mathbf{k}_3} \phi_{-\mathbf{k}_1 - \mathbf{k}_2 - \mathbf{k}_3}.$$

Identify the constants T_c , a_0 , b and u .

- (f) Transform back to real space to obtain the real-space form of $H_{\text{eff}}[\phi_{\mathbf{x}}]$. How is this expression related to Landau theory ?

2. Field theory for the three-state Potts model (10 + 10 + 10 = 30 Punkte)

An extension of the Ising model is the Potts model, which assumes that there are q states at each lattice site and the energy of two sites depends only on whether they are in the same state or not. The Hamiltonian reads

$$H = -\frac{J}{2} \sum_{\langle i, j \rangle} (q\delta_{q_i, q_j} - 1)$$

with $q_i \in \{1, \dots, q\}$ and $\langle i, j \rangle$ denotes the sum over nearest-neighbors. We want to derive the mean-field theory for the $q = 3$ state Potts model using the same prescription as in exercise 1.

- (a) Show that one can represent the energy of this system by $E = \sum_{i, j} J_{ij} \mathbf{s}_i \cdot \mathbf{s}_j$ with \mathbf{s}_i being a two-dimensional normalized vector at site i . Give an explicit representation of these vectors \mathbf{s}_i .
- (b) Following the same prescription as in exercise 1 derive the effective field theory up to quartic order in the Hubbard-Stratonovich fields starting from the microscopic Hamiltonian $H = \sum_{i, j} J_{ij} \mathbf{s}_i \cdot \mathbf{s}_j$ with \mathbf{s}_i being one of the three possible vectors at site i .
- (c) Check to each order in the fields that the terms in $H_{\text{eff}}[\phi_{\mathbf{x}}]$ obey the symmetries of the system.