

Übungen zur Theorie der Kondensierten Materie II SS 16

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Blatt 13

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1. DC Josephson effect

(9 Points)

In this exercise, we will derive the DC Josephson effect microscopically. For this purpose, we consider two identical superconductors described by the BCS mean-field Hamiltonians ($\Delta \in \mathbb{R}$)

$$H_c = \sum_{\mathbf{k}, \sigma} \epsilon_{\mathbf{k}} c_{\mathbf{k}\sigma}^\dagger c_{\mathbf{k}\sigma} - \left(\Delta e^{-i\varphi_c} \sum_{\mathbf{k}} c_{\mathbf{k}\uparrow}^\dagger c_{-\mathbf{k}\downarrow}^\dagger + \text{H.c.} \right), \quad (1a)$$

$$H_d = \sum_{\mathbf{k}, \sigma} \epsilon_{\mathbf{k}} d_{\mathbf{k}\sigma}^\dagger d_{\mathbf{k}\sigma} - \left(\Delta e^{-i\varphi_d} \sum_{\mathbf{k}} d_{\mathbf{k}\uparrow}^\dagger d_{-\mathbf{k}\downarrow}^\dagger + \text{H.c.} \right), \quad (1b)$$

which are coupled by the single-electron-tunneling operator

$$H_t = \sum_{\mathbf{k}, \mathbf{k}', \sigma} t c_{\mathbf{k}\sigma}^\dagger d_{\mathbf{k}'\sigma} + \text{H.c.} \quad (2)$$

The total Hamiltonian of the system is given by $H_c + H_d + H_t$.

- (a) Choose a gauge that removes the phases φ_c and φ_d from Eq. (1). How does the tunneling amplitude t change? (1 Pt.)
- (b) Show that the Josephson current I_J is given by

$$I_J = 2e \langle \partial_{\Delta\varphi} H_t \rangle, \quad \Delta\varphi = \varphi_c - \varphi_d, \quad (3)$$

where e denotes the elementary charge. (2 Pts.)

- (c) Treating H_t as a small perturbation to $H_c + H_d$, expand Eq. (3) to leading nontrivial order in H_t . Perform the contractions using Wick's theorem and bring the result into the form

$$I_J = I_c \sin(\Delta\varphi). \quad (4)$$

Express the critical current I_c in terms of the Nambu Green's functions of a BCS superconductor. (4 Pts.)

- (d) Evaluate the Matsubara and momentum sums to find the explicit form of I_c . You can take the normal state density of states to be independent of energy. (2 Pts.)

2. Spin-susceptibility of a superconductor

(9 Points)

Here we consider an isolated BCS superconductor and calculate its dynamic spin-susceptibility

$$\chi_{jj'}(i\Omega_n, \mathbf{q}) = \int_0^\beta d\tau \langle T_\tau M_j(\mathbf{q}, \tau) M_{j'}(-\mathbf{q}, 0) \rangle e^{i\Omega_n \tau}, \quad M_j(\mathbf{q}, \tau) = \sum_{\mathbf{k}} c_{\mathbf{k}}^\dagger(\tau) \sigma_j c_{\mathbf{k}+\mathbf{q}}(\tau) \quad (5)$$

with $j, j' = 1, 2, 3$, $c_{\mathbf{k}} = (c_{\mathbf{k}\uparrow}, c_{\mathbf{k}\downarrow})^T$, σ_j denoting Pauli matrices, and Ω_n being bosonic. We first work on the imaginary axis and later perform the analytic continuation.

- In a system with spin-rotation symmetry (no spin-orbit coupling), it holds $\chi_{jj'}(i\Omega_n, \mathbf{q}) = \delta_{jj'} \chi_{33}(i\Omega_n, \mathbf{q})$. Why? (1 Pt.)
- Express M_3 in terms of Nambu spinors. Use this to show that the susceptibility can be represented by a bubble diagram where the fermionic lines refer to Nambu Green's functions. (3 Pts.)
- Perform the Matsubara summation to show that

$$\chi_{jj'}(i\Omega_n, \mathbf{q}) = \frac{1}{2} \delta_{jj'} \sum_{\mathbf{k}} \sum_{\lambda, \lambda' = +, -} \left[1 + \frac{\epsilon_{\mathbf{k}} \epsilon_{\mathbf{k}+\mathbf{q}} + \Delta^2}{\lambda \lambda' E_{\mathbf{k}} E_{\mathbf{k}+\mathbf{q}}} \right] \frac{n_F(\lambda' E_{\mathbf{k}+\mathbf{q}}) - n_F(\lambda E_{\mathbf{k}})}{i\Omega_n - (\lambda' E_{\mathbf{k}+\mathbf{q}} - \lambda E_{\mathbf{k}})}, \quad (6)$$

where $E_{\mathbf{k}} = \sqrt{\epsilon_{\mathbf{k}}^2 + \Delta^2}$ and n_F is the Fermi function. (2 Pts.)

- For example, the NMR relaxation rate is determined by

$$\frac{1}{T_1 T} \propto \sum_{\mathbf{q}} \lim_{\Omega \rightarrow 0} \frac{\text{Im} \chi_{jj}^R(\Omega, \mathbf{q})}{\Omega}. \quad (7)$$

Using the expression (6) of the susceptibility, the analytic continuation to the retarded function $\chi_{jj}^R(\Omega, \mathbf{q})$ is straightforward. Write the right hand side of Eq. (7) as a single integral over the quasi-particle energies E of the superconductor (the normal state density of states can be approximated as a constant). (3 Pts.)