

Dr. Boris Narozhny, Institut für Theorie der Kondensierten Materie (TKM), narozhny@tkm.uni-karlsruhe.de
 Dr. Panagiotis Kotetes, Institut für Theoretische Festkörperphysik (TFP), panagiotis.kotetes@kit.edu

THEORETICAL OPTICS: EXERCISE SHEET 3

Announcement date: 05.05.2013 – Tutorials' date: 09.05.2013 and 10.05.2013

1. Wave propagation in optically anisotropic materials - Fresnel's equation

For optically anisotropic and magnetically non-active lossless media, we have the constitutive relations $\mathbf{D}(\mathbf{k}, \omega) = \varepsilon_0 \boldsymbol{\varepsilon}(\omega) \boldsymbol{\mathcal{E}}(\mathbf{k}, \omega)$ and $\mathbf{B}(\mathbf{k}, \omega) = \mu_0 \mathbf{H}(\mathbf{k}, \omega)$. For these materials, the dielectric constant $\varepsilon(\omega)$ is replaced by the dielectric tensor $\boldsymbol{\varepsilon}(\omega)$. For the particular case we consider

$$\boldsymbol{\varepsilon}(\omega) = \begin{pmatrix} \varepsilon_{xx}(\omega) & 0 & 0 \\ 0 & \varepsilon_{yy}(\omega) & 0 \\ 0 & 0 & \varepsilon_{zz}(\omega) \end{pmatrix}. \quad (1)$$

a. Combine Eq. (1) with

$$\mathbf{D}(\mathbf{k}, \omega) = \varepsilon_0 (ck/\omega)^2 \left\{ \boldsymbol{\mathcal{E}}(\mathbf{k}, \omega) - [\hat{\mathbf{k}} \cdot \boldsymbol{\mathcal{E}}(\mathbf{k}, \omega)] \hat{\mathbf{k}} \right\}, \quad (2)$$

where $\hat{\mathbf{k}} = \mathbf{k}/|\mathbf{k}|$ and $k = |\mathbf{k}|$, in order to obtain a homogeneous system of equations for the components of $\boldsymbol{\mathcal{E}}$. For non-zero solutions of $\boldsymbol{\mathcal{E}}$, derive the so called Fresnel's equation, which provides the corresponding dispersion relations $\omega(\mathbf{k})$. (3 points)

b. Derive the dispersion relations $\omega(\mathbf{k})$ for an optically uni-axial material with $\varepsilon_{xx}(\omega) = \varepsilon_{yy}(\omega) \equiv \varepsilon_{\perp}$ and $\varepsilon_{zz}(\omega) = \varepsilon_{\parallel}$, where the z -axis defines the optic axis. (3 points)

c. Calculate the group velocity vector $\mathbf{v}_g = \frac{d\omega(\mathbf{k})}{d\mathbf{k}}$, for the dispersion relations found in 1.b. (2 points)

2. Electromagnetic energy flux density in optically anisotropic materials

Consider the following electromagnetic wave

$$\boldsymbol{\mathcal{E}}(\mathbf{r}, t) = \mathcal{E}_0 \hat{\mathbf{n}} \cos[\mathbf{k} \cdot \mathbf{r} - \omega_e(\mathbf{k})t], \quad (3)$$

propagating in an optically uni-axial material with $\varepsilon_{xx}(\omega) = \varepsilon_{yy}(\omega) \equiv \varepsilon_{\perp}$ and $\varepsilon_{zz}(\omega) = \varepsilon_{\parallel}$. The wave propagates with wave-vector $\mathbf{k} = \frac{k_0}{\sqrt{2}}(0, 1, 1)$ according to the extraordinary (e) dispersion relation

$$\omega_e(\mathbf{k}) = c \sqrt{\frac{\varepsilon_{\perp} k_{\perp}^2 + \varepsilon_{\parallel} k_{\parallel}^2}{\varepsilon_{\parallel} \varepsilon_{\perp}}} \quad \text{where} \quad k_{\perp} = \sqrt{k_x^2 + k_y^2} \quad \text{and} \quad k_{\parallel} = k_z. \quad (4)$$

a. Determine the polarization of the electric field, $\hat{\mathbf{n}}$, based on Eq.(2). (4 points)

b. Calculate the average energy flux density

$$\langle \mathbf{S} \rangle_T = \langle \boldsymbol{\mathcal{E}}(\mathbf{r}, t) \times \mathbf{H}(\mathbf{r}, t) \rangle_T = \frac{1}{T} \int_0^T \boldsymbol{\mathcal{E}}(\mathbf{r}, t) \times \mathbf{H}(\mathbf{r}, t) dt, \quad (5)$$

over a period $T = 2\pi/\omega_e(\mathbf{k})$, for the electromagnetic wave of 2.a. (4 points)

3. Refraction of electromagnetic waves in optically anisotropic materials

Consider an electromagnetic wave incident to a surface separating vacuum from an optically uni-axial material with $\varepsilon_{xx}(\omega) = \varepsilon_{yy}(\omega) \equiv \varepsilon_{\perp}$ and $\varepsilon_{zz}(\omega) = \varepsilon_{\parallel}$. For zero surface charge and current densities, the boundary conditions for \mathbf{D} , \mathcal{E} , \mathcal{B} and \mathbf{H} read

$$\hat{\mathbf{n}} \cdot (\mathbf{D}_1 - \mathbf{D}_2) = 0, \quad \hat{\mathbf{n}} \cdot (\mathcal{B}_1 - \mathcal{B}_2) = 0, \quad \hat{\mathbf{n}} \times (\mathcal{E}_1 - \mathcal{E}_2) = 0 \quad \text{and} \quad \hat{\mathbf{n}} \times (\mathbf{H}_1 - \mathbf{H}_2) = 0, \quad (6)$$

where $\hat{\mathbf{n}}$ is the unit vector perpendicular to the surface and 1,2 denote the corresponding fields in vacuum and the dielectric respectively.

- Using the above boundary conditions, retrieve the expressions for the reflected \mathcal{E}_r and transmitted \mathcal{E}_t electric fields. (10 points)
- Draw the wave-vector surfaces $\omega(\mathbf{k}) = \text{constant}$ for the allowed dispersion relations in the dielectric. Consider both possibilities: **i.** positive uni-axial materials with $\varepsilon_{\parallel} > \varepsilon_{\perp}$ (for instance quartz SiO_2) and **ii.** negative uni-axial materials with $\varepsilon_{\parallel} < \varepsilon_{\perp}$ (for instance calcite CaCO_3). (4 points)

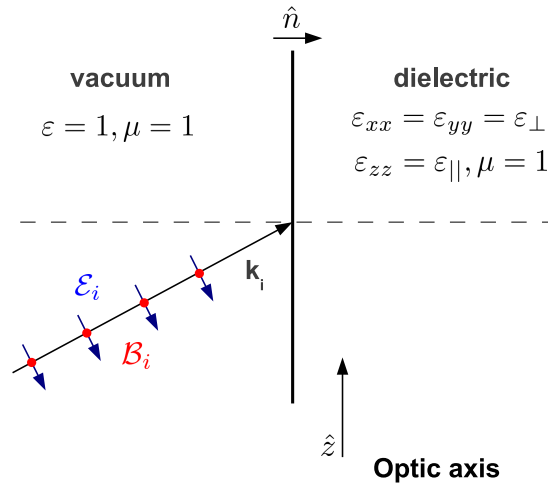


FIG. 1: Refraction of an electromagnetic wave \mathcal{E}_i incident to the surface separating vacuum from an optically active uni-axial dielectric. The polarization of the electric field, \mathbf{k}_i and the optic axis lie within the same plane.